The 2024 International Workshop on Future Linear Colliders, LCWS2024[1], took place on July 8-11, 2024 at the University of Tokyo, Japan. The LCWS workshop series is devoted to the study of the physics, detectors, and accelerator issues relating to high-energy linear electron-positron colliders. A linear collider will initially operate as a Higgs factory, as well as providing a clear path for upgrades in both energy and luminosity.

The 2024 workshop had 341 registered participants, 254 of who joined in-person and 87 by remote connection, and a total of around 290 contributions. The workshop consisted of a mix of plenary and parallel sessions. Parallel sessions related to physics and detectors were grouped into:

- Higgs and Electroweak physics
- Top quark, QCD, and Flavor physics, and Precision Modelling
- BSM physics and Global Interpretations
- Vertex, Tracking, and Timing detectors
- Calorimetry and Muon detectors
- Software, Reconstruction, and Computing,

while accelerator-related sessions covered

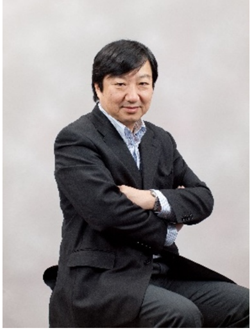
- Sources
- Damping Rings and Beam Delivery Systems
- Beam Dynamics
- Super-conducting RF
- Normal-conducting RF
- Advanced Accelerator Concepts
- Applications
- Conventional Facilities and Machine Detector interface.

There were additional sessions devoted to Sustainability, Industry, and Early Career Researchers, and a lively poster session featuring 18 posters mostly by young researchers and students, four of whom were awarded prizes.

A special session was devoted to the memory of Sachio Komamiya, a long-standing supporter of linear colliders, who had died a few weeks before the workshop. He played a leading role in the advancement and promotion of the International Linear Collider (ILC) project, both internationally and within the Japanese high energy physics community. He served

[1] https://agenda.linearcollider.org/e/lcws2024

multiple terms as chairperson of the High Energy Physics Committee of JAHEP (Japan Association of High Energy Physicists), and as chairperson of the Linear Collider Board during 2013-2017. Many happy memories of Sachio were recounted by his friends and colleagues from around the world.



Sachio Komamiya

The global situation around a future Higgs Factory is evolving rapidly. The feasibility study of the FCC-ee, a possible future project for CERN, is reaching its conclusion, and the CEPC project is preparing to submit a proposal to the Chinese government in the next year. The process to update the European Strategy in Particle Physics is underway. In a discussion session on the "Global vision for a Linear Collider facility", LCWS2024 participants agreed that a strong proposal should be made to the European Strategy for a Linear Collider as the next large accelerator facility, emphasising both its initial role as a Higgs Factory and the various technological paths which can be used to upgrade such a collider to the TeV scale.

The LCWS2024 was sponsored by a number of organisations and companies, whose support was essential in ensuring a fruitful and enjoyable workshop. We would also like to thank members of the International Advisory and Program Commitees, Session Conveners, and all staff, contributors and attendees for making such an interesting and lively workshop.

Taikan Suehara (ICEPP, U. Tokyo)
Chair of the Local Organizing Committee

Junping Tian (ICEPP, U. Tokyo), *Daniel* Jeans (KEK)
Editors of these Proceedings

December, 2024

LCWS2024 participants

Heavy top quark mass in the minimal universal seesaw model

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Abstract. We study the hierarchy between M_T, v_L, and v_R, the relevant energy scales of the Minimal Universal Seesaw Model (MUSM), where the two lightest quark families remain massless at tree level. We also predict the heavy top quark mass, $m_{t'}$. We do some numerical analysis using recent experimental data. Our numerical analysis demonstrates that M_T is sensitive to the values of the Yukawa couplings. The heavy top quark mass ($m_{t'}$) is predicted to be within the range from 1.4 TeV to 7.2 TeV for $v_R = 10$ TeV.

1 Introduction

One of the questions in the particle physics is the origin of the fermion mass hierarchy. Fermion masses vary widely, with no clear explanation for the large differences between them. This hierarchy is particularly evident in the quark sector. The discovery of the top quark by the CDF [1] and DØ [2] collaborations in 1995 further highlighted this issue. The observed top quark is much heavier than other quarks, which presents a puzzle in understanding the origin of mass differences in the quark sector. This mass hierarchy problem has motivated to extend the SM. One extension of the SM that attempts to explain the fermion mass hierarchy is the Universal Seesaw Model (USM) prior top quark discovery in Refs. [3–9] and after top quark discovery, e.g., in Refs. [10–12].

A more recent development of the USM, where the two lightest quark families are massless at tree level, has been proposed in Ref. [13]. This model, called the Minimal Universal Seesaw Model (MUSM), naturally explains the observed quark mass hierarchy in the third family. Some of its phenomenological implications, such as flavor-changing neutral current (FCNC) processes in the interactions between the Higgs and Z bosons with quarks, are discussed in detail. One of the new physics (NP) particles predicted by this model is the heavy top quark (t'), also referred to as a top quark partner in some references. Searches for this heavy top quark have been conducted, e.g., by the CMS [15] and ATLAS [16] collaborations and have been summarized by the Particle Data Group (PDG) [17]. These results provide a lower bound for the mass of the heavy top quark. In this work, we will study the prediction of the heavy top quark mass in more detail in Ref. [13].

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Table 1. Particle contents of the model, where $i \in \{1, 2, 3\}$ is the family index. The content of this table is taken from Table 1 in Ref.[13].

Quark and Higgs fields	$SU(3)_C$	$SU(2)_L$	$SU(2)_R$	$U(1)_{Y'}$
$q^i_L = \begin{pmatrix} u^i_L \\ d^i_L \end{pmatrix}$	**3**	**2**	**1**	1/6
$q^i_R = \begin{pmatrix} u^i_R \\ d^i_R \end{pmatrix}$	**3**	**1**	**2**	1/6
$T_{L,R}$	**3**	**1**	**1**	2/3
$B_{L,R}$	**3**	**1**	**1**	−1/3
$\phi_L = \begin{pmatrix} \chi^+_L \\ \chi^0_L \end{pmatrix}$	**1**	**2**	**1**	1/2
$\phi_R = \begin{pmatrix} \chi^+_R \\ \chi^0_R \end{pmatrix}$	**1**	**1**	**2**	1/2

2 The model

In this section, we briefly introduce the model in Ref. [13]. The model is based on $SU(3)_C \times SU(2)_L \times SU(2)_R \times U(1)_{Y'}$ gauge symmetry. The particle content of the model is SM particle with additional one up-type (T) and one down-type (B) vector-like quarks (VLQs) as partner of top and bottom quark, respectively. Additionally, there is $SU(2)_R$ Higgs doublet (ϕ_R). The charge convention that used in the model is

$$Q = I^3_L + I^3_R + Y', \tag{1}$$

where Q, $I^3_{L(R)}$, and Y' are electromagnetic charge, left(right) weak-isospin, and $U(1)_{Y'}$ hypercharge, respectively. The particle contents of the model are given in table 1.

The Lagrangian of Yukawa interactions and mass terms of VLQs is as follows [13]:

$$\begin{aligned} \mathcal{L}_{\text{yuk}} = & -Y^3_{u_L}\overline{q^3_L}\tilde{\phi}_L T_R - Y^3_{u_R}\overline{T_L}\tilde{\phi}^\dagger_R q^3_R - \overline{q^i_L}y^i_{d_L}\phi_L B_R - \overline{B_L}y^{i*}_{d_R}\phi^\dagger_R q^i_R - h.c. \\ & - \overline{T_L}M_T T_R - \overline{B_L}M_B B_R - h.c., \end{aligned} \tag{2}$$

where the Yukawa couplings of the up-type quark ($Y^3_{u_L}$ and $Y^3_{u_R}$) are real positive numbers in a specific weak-basis while the Yukawa couplings of the down-type quark ($y^i_{d_L}$ and $y^{i*}_{d_R}$) are general complex vectors[1]. The charge conjugation of Higgs fields is defined as $\tilde{\phi}_{L(R)} = i\tau^2\phi^*_{L(R)}$ where τ^a with $a \in \{1, 2, 3\}$ is the Pauli matrix. In the second line of Eq.(2), M_T and M_B are the VLQ mass parameters that were taken as real numbers.

The symmetry of the model is spontaneously broken in two stages. First, $SU(2)_R \times U(1)_{Y'} \to U(1)_Y$ when the neutral component of the $SU(2)_R$ Higgs doublet acquires non-zero vacuum expectation value (vev) v_R. After this first stage of symmetry breaking, the symmetry of the model reduces to the SM gauge symmetry, $SU(3)_C \times SU(2)_L \times U(1)_Y$. In the second stage, $SU(2)_L \times U(1)_Y \to U(1)_{em}$ when the neutral component of the $SU(2)_L$ Higgs doublet acquires non-zero vev v_L. The definition of the vevs are given as follows:

$$\langle\phi_R\rangle = \frac{1}{\sqrt{2}}\begin{pmatrix} 0 \\ v_R \end{pmatrix}, \qquad \langle\phi_L\rangle = \frac{1}{\sqrt{2}}\begin{pmatrix} 0 \\ v_L \end{pmatrix} \tag{3}$$

[1]The details are given in Appendix A of Ref. [13].

and satisfy $v_R \gg v_L$. One can follow the derivation in Ref. [13] and obtain the exact mass eigenvalue of top quark (t) and heavy top quark (t') as follows[2]:

$$m_t = \frac{\sqrt{M_T^2 + (m_{u_R} + m_{u_L})^2}}{2} - \frac{\sqrt{M_T^2 + (m_{u_R} - m_{u_L})^2}}{2}, \tag{4}$$

$$m_{t'} = \frac{\sqrt{M_T^2 + (m_{u_R} + m_{u_L})^2}}{2} + \frac{\sqrt{M_T^2 + (m_{u_R} - m_{u_L})^2}}{2}, \tag{5}$$

where,

$$m_{u_R} = Y^3_{u_R} \frac{v_R}{\sqrt{2}}, \qquad m_{u_L} = Y^3_{u_L} \frac{v_L}{\sqrt{2}}. \tag{6}$$

3 Numerical analysis

In this section, we discuss the hierarchy between M_T, v_L, and v_R using the exact mass eigenvalues of the top quark and the heavy top quark given in Eqs. (4) and (5), respectively. We analyze the constraints on M_T and v_R, while taking $v_L = 246.22$ GeV from Ref. [17]. As explained in Ref. [13], the lower bound constraint on v_R is $v_R \gtrsim 10$ TeV[3]. Additionally, we use numerical input from Ref. [17] where the top quark mass is $m_t = 172.57$ GeV and the lower bound for heavy top quark mass is $m_{t'} > 1310$ GeV. The Yukawa couplings $Y^3_{u_R}$ and $Y^3_{u_L}$ are free parameters. However, we take the upper limit of the Yukawa couplings are $Y^3_{u_R}, Y^3_{u_L} \leq 1$.

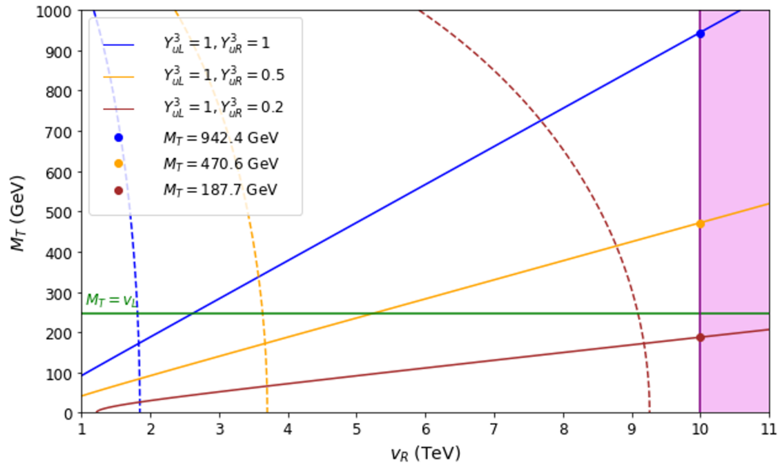

Figure 1. The variation of M_T as a function of v_R, for different values of $Y^3_{u_R}$, with $Y^3_{u_L} = 1$. The vertical purple line represents the lower limit of $v_R = 10$ TeV. The solid blue, orange, and brown lines represent the constraints from the exact mass eigenvalue of the top quark ($m_t = 172.57$ GeV) for $Y^3_{u_R} = 1, 0.5, 0.2$, respectively. The corresponding dashed lines show the lower limit of the heavy top quark mass ($m_{t'} > 1310$ GeV) for each value of $Y^3_{u_R}$. The dots on the lines at $v_R = 10$ TeV indicate the respective M_T values. The shaded pink region represents the allowed parameter space for $v_R \geq 10$ TeV.

[2] This exact mass eigenvalue was first introduced in Ref. [14]. Details on the diagonalization of the quark mass matrix can be found in Appendix C of Ref. [13]. In this work, we focus solely on the top sector.

[3] This lower bound is derived from the lower limit on the mass of the Z' boson, which in turn constrains the mass of the W_R boson mass. Assuming $g_R \simeq 1$, this gives the lower bound on v_R. For more details, see Ref. [13].

Table 2. The heavy top quark mass for varying values of the $Y^3_{u_R}$ while keeping $Y^3_{u_L} = 1$.

$Y^3_{u_R}$	M_T	$m_{t'}$
1	942.4 GeV	7.13 TeV
0.5	470.6 GeV	3.57 TeV
0.2	187.7 GeV	1.43 TeV

In figure 1, we explore the dependence of M_T on v_R for varying values of the $Y^3_{u_R}$ while keeping $Y^3_{u_L} = 1$. The plot illustrates that as $Y^3_{u_R}$ decreases, the corresponding M_T values also decrease for a given v_R. However, the variation of $Y^3_{u_L}$ is more stringently limited. It must satisfy $Y^3_{u_L} \geq y_t^{\text{SM}}$ [13], where y_t^{SM} is the SM Yukawa coupling of the top quark, with a value of $y_t^{\text{SM}} \simeq 0.9912$.

Another important result is that two possible hierarchies emerge: $v_L < M_T < v_R$ or $M_T < v_L < v_R$, depending on the chosen parameters. After obtaining M_T, we compute the heavy top quark mass ($m_{t'}$) using Eq. (5) and summarize the results in table 2. The predicted heavy top quark mass is in the range of approximately 1.4 to 7.2 TeV for $v_R = 10$ TeV, which could be tested in future linear collider experiments.

4 Conclusion

We have investigated the hierarchy between M_T, v_L, and v_R, and predicted the heavy top quark mass within the Minimal Universal Seesaw Model (MUSM), where the two lightest quark families remain massless at tree level. Our numerical analysis demonstrates that M_T is highly sensitive to the values of the Yukawa couplings. Most importantly, the model predicts a heavy top quark mass in the range of approximately 1.4 to 7.2 TeV for $v_R = 10$ TeV. This mass range represents a significant prediction that could be tested in future collider experiments, such as those planned for future linear colliders.

Acknowledgements

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